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ELASTIC STRESSES AT A MISMATCHED CIRCUMFERENTIAL JOINT IN A PRESSURIZED CYLINDER INCLUDING THICKNESS CHANGES AND MERIDIONAL LOAD COUPLING

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SUMMARY

The results of an analytical investigation of stresses at a circumferential joint in a pressurized circular cylindrical vessel are presented graphically for a wide range of geometric values. This investigation of a mismatched joint with any thickness change was performed for three types of junction geometry: unfilleted butt joint, filleted butt joint, and overlap joint. These analyses include nonlinear effects for large radius-to-thickness structures that are highly stressed. This report provides design curves that predict the stresses at such junctions; also illustrated are the conditions under which the meridional load coupling effect becomes important.

All thickness ratios for the three joints were investigated. In the butt joints, the mismatch was varied from zero to the case where the inner surface of one cylinder coincided with the outer surface of the other cylinder. The lengths of overlap joints investigated varied from short lengths with discontinuity interactions to longer lengths with no interactions. Nonlinear edge influence coefficients, which were used to obtain the results, are presented for both semi-infinite and short-length cylinders.

It was concluded that, for a change in thickness at a circumferential joint in a cylindrical pressure vessel, it is less detrimental if the thinner section is displaced a small amount inward rather than outward from the thicker section if mismatch is to occur. It was also concluded that, if the thickness does not change at the junction of an unfilleted butt joint, a nonlinear stress analysis is not required to determine the mismatch stresses. For equal thicknesses at a filleted butt joint, a nonlinear analysis is required to determine accurately the mismatch stresses.

INTRODUCTION

Cylindrical shells that are fabricated in sections and then joined to form the completed structure are typical of many propellant tanks and pressure vessels. The sections joined circumferentially are sometimes of different thicknesses, and mismatch is often introduced during the joining process. It may be necessary to consider the effect of the meridional load on the bending stresses in order to determine accurately the stresses for structures that are highly stressed and have a large radius-to-thickness ratio. This consideration results in a nonlinear analysis in which stresses are not proportional to internal pressure.

Previous investigations of related subjects did not include all the parameters considered herein and were usually linear in nature. The analyses presented in references 1 and 2 are linear. Reference 3 is applicable only for mismatch at junctions of equal thickness but in any stress field. References 4 to 6 are nonlinear analyses that include thickness changes but not mismatch. Reference 7 is an outline of a computer program for the analysis of reference 6. Reference 8, which includes mismatch effects, is an extension of reference 5. All these analyses are for unfilleted butt joints only; no considerations of filleted butt joints or overlap joints were included. An analysis of a junction of equal thickness cylinders by an overlap of short length is given in reference 9. Excellent agreement between linear theory and experimental results for small scale tanks is presented in references 10 and 11.

The nonlinear analyses of this investigation considered thickness changes and mismatch for unfilleted and filleted butt joints as well as for overlap joints. This report provides a compilation of design curves for determining the maximum stresses near circumferential joints in cylindrical pressure vessels considering nonlinear effects. The conditions under which the nonlinear effects become important and necessitate consideration are also illustrated.

An analytical investigation was performed to determine accurately the stresses near three types of circumferential joints: unfilleted butt, filleted butt, and overlap. Results for these three joints are presented graphically for any thickness ratio. In the butt joints, the mismatch studied varied from zero to the case where the inner surface of one component coincided with the outer surface of the other component. The length of overlap joints investigated varied from short lengths with interactions to longer lengths with no interactions.

GEOMETRY AND METHOD OF ANALYSIS

General

Two semi-infinite closed-end circular cylinders (fig. 1(a)), which are joined in one

of the three ways shown in figure 1(b), are considered. As can be seen, both the filleted butt joint and the overlap joint have two discontinuities. For these circumferential joints, the analysis assumes interactions between the two discontinuities of the joint. The cylinders are loaded by internal pressure only. It is assumed that the two cylinders are coaxial and that the radius and thickness of each are constant; however, a mismatch (eccentricity) between the middle surfaces at the joint is considered. Both cylinders are made of the same material, which is assumed to be homogeneous, isotropic, and subject to Hooke's law. The effect of stress concentration due to sharp corners at the discontinuity is not included.

Unfilleted Butt Joint

The geometry of an unfilleted butt joint (fig. 2(a)) can be described by the thickness ratio c and the percent mismatch m. (All symbols are defined in appendix A.) If the subscript 1 is assigned to the smaller thickness component and the subscript 2 to the larger thickness component, the thickness ratio can be defined as

$$c = \frac{h_1}{h_2} \tag{1}$$

When the same subscripts are kept, the percent mismatch is defined as

$$m = \frac{a_2 - a_1}{(h_1 + h_2)/2} (100)$$
 (2)

It should be noted that the percent mismatch can be positive or negative depending on whether $a_2 - a_1$ is positive or negative.

The method used to determine the stresses at the joint is the usual edge influence coefficient approach in which the cylinder is imagined cut at the discontinuity and separated into two edge-loaded cylinders. Simultaneous solution of the equations that express the equality of deflection and rotation at the discontinuity gives the shear forces and bending moments at the junction. The shear forces and bending moments for an unfilleted butt joint are given in appendix B.

The necessary edge influence coefficients were obtained from reference 12 by analogy to a semi-infinite beam on an elastic foundation with a free end. For a nonlinear analysis, the beam is loaded simultaneously with a unit end shear force or bending

moment and an axial tensile force. Nonlinear edge influence coefficients are presented in equation (B5).

The nonlinear influence coefficients differ from the usual linear influence coefficients by quantities that contain a nonlinearity parameter n, which is defined as

$$n = \frac{\sqrt{3(1 - \nu^2)}}{2} \frac{p}{E} \left(\frac{a}{h}\right)^2$$
 (3)

The parameter n may be determined from figure 3. It is emphasized that this parameter is a function of internal pressure as well as geometry. If n is set equal to zero, the nonlinear analysis reduces to the linear analysis.

Filleted Butt Joint

The filleted butt joint was assumed to have a 45° fillet, as shown in figure 2(b). The analysis for the filleted butt joint is similar to that of the unfilleted butt joint, except that the fillet area was treated as a ring that joins the two cylinders. The structure was imagined to be cut at each discontinuity and separated into three components. The ring cross section was assumed to translate and rotate but not deform. All loads on the circular ring were resolved to radial forces and twisting couples uniformly distributed along the centerline. The influence coefficients were then determined by modification of the familiar ring formulas that are given in references 13 and 14.

The equality of deflections and rotations at each of the two discontinuities of the filleted butt joint gave the four equations required to determine the unknown shear forces and bending moments at both discontinuities. These shear forces and bending moments for a filleted butt joint are given in appendix C.

Overlap Joint

The overlap joint is shown in figure 2(c). The analysis for this joint was identical to that for the filleted butt joint except that the center component was analyzed as a short cylinder instead of as a ring. The edge influence coefficients were taken from reference 12 for a finite-length beam on an elastic foundation.

The equality of deflections and rotations at each of the two discontinuities of the overlap joint gave the four equations required to determine the unknown shear forces and bending moments at both discontinuities. These shear forces and bending moments for

an overlap joint are given in appendix D.

In the overlap joint, the percent mismatch is not applicable, since all joints have positive or negative 100-percent mismatch. The quantity required to describe an overlap joint, in addition to thickness ratio and nonlinearity parameter, is the overlap parameter $\lambda_3 \ell$, which is defined as

$$\lambda_3 \ell = \sqrt{\frac{3(1 - \nu^2)}{a_3^2 h_3^2}} \ell \tag{4}$$

where λ_3 is the characteristic of the short center cylindrical section and ℓ the length of overlap parameter is dimensionless.

Stress Determination

The total stress at any location is composed of membrane stresses and discontinuity stresses. The stress distribution was not investigated in the ring or short cylinder components because the stresses were relatively low due to the larger thicknesses. Equations for the stress distribution in the cylindrical components are given in appendix E.

The energy of distortion theory, introduced by von Mises and reinterpreted by Hencky, can be used to predict the limit of elastic behavior. This theory states that yielding will occur in a biaxial state of stress if a quantity called the effective stress is equal to the yield stress as determined from a uniaxial tensile test. The effective stress is defined as

$$\sigma_{\mathbf{e}} = \sqrt{\sigma_{\theta}^2 + \sigma_{\mathbf{x}}^2 - \sigma_{\theta}\sigma_{\mathbf{x}}} \tag{5}$$

The membrane stresses in the thinner cylinder are used to determine the maximum effective membrane stress. The maximum effective stress anywhere in the structure is then determined by using the total stresses. The ratio of the maximum effective stress anywhere in the vessel to the maximum effective membrane stress is defined herein as the stress factor K.

SUMMARY OF ASSUMPTIONS AND LIMITATIONS

The following assumptions and limitations apply to the analyses performed herein

and to the application of the results.

- (1) Small finite deflection theory is assumed. (Normal deflections of the order of the thickness are permissible.)
 - (2) Thin shell theory is applicable.
 - (3) The cylinder has closed ends and is loaded by uniform internal pressure.
- (4) It is assumed that the thickness of each circular cylindrical component is constant and that the axes of revolution are collinear.
 - (5) The material is assumed to be homogeneous, isotropic, and linearly elastic.
- (6) The effect of stress concentrations at sharp corners is not included (i.e., plane sections remain plane).
- (7) For the filleted butt joint, the analysis assumes that the ring radius is much larger than its thickness and that the ring has a rigid cross section.
 - (8) For the overlap joint, short cylinder equations are used in the overlap portion.
- (9) Stress factors are based on the effective stress as defined in the energy of distortion theory of yielding.

RESULTS AND DISCUSSION

Influence Coefficients

Nonlinear edge-rotation and edge-deflection influence coefficients due to edge shear force and bending moment on a circular cylinder are presented in equation (B5). These influence coefficients can be used for other shapes of shells in accordance with Geckeler's approximation (ref. 15, pp. 567-8). With these influence coefficients, nonlinear results can be obtained for discontinuity stresses at the junctions of shells of other geometry in which the edge tangents are parallel to the axis of revolution.

Design Curves

For all three joints, the complexity and length of the equations did not readily permit reduction to simple form. Consequently, values for geometry (a, h), material constants (E, ν) , and internal pressure (p) were selected for use in the computer programs to obtain the stress factors (K).

All results are presented for a Poisson's ratio of 1/3. A comparison with a Poisson's ratio of 0.30 resulted in a maximum stress factor variation of approximately 1.2 percent.

Values of the nonlinearity parameter factors a, h, p, and E were varied, but the value of the nonlinearity parameter itself was unchanged; a comparison of the corresponding

results showed no appreciable change in K. It was concluded that K in the nonlinear range is a function of n and not its individual factors.

Stress factors for the three joints are presented for all thickness ratios. Each joint was investigated for n parameters of 0, 0.1, 1, 10, and 100. For all curves presented herein, solid lines represent configurations in which the maximum effective stress is located at a discontinuity, and the dashed lines represent configurations in which the maximum effective stress location is away from a discontinuity.

<u>Unfilleted butt joint</u>. - The results for the unfilleted butt joint are presented graphically in figure 4. Stress factors K are presented for mismatch from negative to positive 100 percent.

Negative mismatch (thinner cylinder displaced outward from thicker cylinder) is almost always more critical than positive mismatch of the same magnitude. This is especially true for the smaller thickness ratios, where the thicknesses are quite dissimilar. When $n \ge 1.0$, K = 1.0 if the maximum effective stress location is not at the discontinuity; thus, there is no increase in K because of the discontinuity in geometry.

Filleted butt joint. - All configurations investigated with an unfilleted butt joint were also investigated with a 45° filleted butt joint, as shown in figure 2(b). Stress factors were determined for mismatch from negative to positive 100 percent for all thickness ratios. The stress factors for the filleted butt joint are presented in figure 5.

The curves in figure 5 are similar to the corresponding unfilleted butt joint curves in figure 4. It should be noted that, as the nonlinearity parameter increases, the stress factor is reduced for most cases. For equal thicknesses, up to a 25-percent decrease in the stress factor is possible. As for the unfilleted butt joint, when $n \ge 1.0$, K = 1.0 if the maximum effective stress is not at the discontinuity; thus, there is no increase in K because of the discontinuity in geometry.

Overlap joint. - The overlap joint was investigated for overlap parameters $\lambda_3 \ell$ from 0 to 10 for all thickness ratios. Stress factors for these cases are presented in figure 6. Results for the overlap parameter of zero were taken from corresponding cases of a positive or negative 100-percent unfilleted butt joint.

When the overlap parameter is 10 or greater, no interactions occur between the discontinuities of the joint; the junction can be treated, therefore, as two unfilleted butt joints that are far enough apart that they do not interact. The stress factors for overlap joints with an overlap parameter of 10 were compared with the corresponding results for an unfilleted butt joint; the results are identical.

As can be seen in figure 6, overlap joints with positive overlap are usually preferred to corresponding joints with negative overlap; that is, it is preferable from a stress viewpoint to mate the outside surface of the thinner component to the inside surface of the thicker component. For negative overlap, the stress factor decreases as the overlap parameter is increased. For positive overlap, however, a long overlap length is desired

only for thickness ratios greater than approximately 0.25. For smaller thickness ratios, a shorter overlap length is preferred so that the lowest stress factors possible can be obtained. For negative overlap and a given overlap parameter, the stress factor decreases with increasing nonlinearity. For positive overlap and a given overlap length, the stress factor increases for some thickness ratios but decreases for equal thicknesses as well as for smaller thickness ratios with increasing nonlinearity. Where $n \ge 1.0$, K = 1.0 if the maximum effective stress is not at the discontinuity.

Comparisons of Results

Pressurized cylinder with fixed ends. - An analysis was performed for a cylindrical shell that had fixed ends (with longitudinal expansion possible) and that was loaded by a uniform internal pressure. A nonlinear analysis for the fixed-end cylinder resulted in an equation for a stress factor that was independent of the nonlinearity parameter n. It was found that this equation, which is a function of Poisson's ratio only, is the same as would be obtained from the usual linear analysis. These results are given in figure 7.

From the results shown in figure 7 for a fixed-end cylinder with $\nu=1/3$, a value of K=2.068 is obtained. It is observed in figure 8 that the fixed-end cylinder is approached for very small thickness ratios, and that, for these thickness ratios, K is linear with percent mismatch.

Filleted compared with unfilleted butt joint. - A comparison of the stress factors for unfilleted butt joints (fig. 4) with corresponding filleted butt joints (fig. 5) was made to determine the effect on K of adding a 45° fillet. It is emphasized that stress concentration at sharp corners has been neglected in this comparison. For equal thicknesses at the discontinuity, linear results show a 2-percent decrease in the stress factor when the 45° fillet is added; the decrease did not exceed 4 percent for any thickness ratio.

For the nonlinear results, a considerable decrease in the stress factor is obtained by adding the 45^{O} fillet. For n = 100, this decrease can be as much as 25 percent for most thickness ratios (including equal thicknesses).

Linear compared with nonlinear results. - One result of a comparison of linear and nonlinear results for unfilleted butt joints with positive mismatch and thickness ratios greater than about 0.17 but less than 0.8 is that a linear analysis may be unconservative and give a stress factor as much as 30 percent too low. For these joints with positive mismatch, the stress factor K increases as the nonlinearity parameter n increases. For most other unfilleted butt joints, the stress factor decreases with increasing nonlinearity.

Another result of a comparison between linear and corresponding nonlinear unfilleted butt joints was that, for equal thicknesses at the joint, c = 1.0, the stress factors were independent of the amount of nonlinearity. When this comparison was performed for a

filleted butt joint with equal thicknesses, a considerable decrease (up to 25 percent) in stress factor with increasing nonlinearity was noted.

For the three joints investigated, if the nonlinearity parameter n was about 0.8 or less, the difference between corresponding linear and nonlinear cases was no more than 5 percent. Thus, for n less than about 0.8, linear junction analyses can be used.

SAMPLE PROBLEM

A sample problem illustrates the use of the design curves. The configuration selected for the sample problem is illustrated in figure 9. Two stainless-steel cylinders of 120-inch nominal diameter have a 1/2-inch overlap and are pressurized to 35 pounds per square inch.

The necessary constants are the following:

Middle surface radius of thinner cylinder, a_1 , in.			•					60
Middle surface radius of thicker cylinder, a ₂ , in.							6	0.0135
Thickness of shell wall of thinner cylinder, h_1 , in.								
Thickness of shell wall of thicker cylinder, h ₂ , in.								0.014
Length of overlap, ℓ , in								. 0.5
Poisson's ratio, ν								. 0.3
Modulus of elasticity, E, psi		. :.						30×10 ⁶
Internal pressure, p, psi								

Since the thicker cylinder is displaced outward from the thinner cylinder, this is positive overlap. To completely specify a positive overlap joint, it is necessary to determine the thickness ratio, overlap parameter, and nonlinearity parameter:

Thickness ratio:

$$c = \frac{h_1}{h_2} = \frac{0.013}{0.014} = 0.929$$

Overlap parameter:

$$\lambda_{3}\ell = \sqrt[4]{\frac{3(1-\nu^{2})}{a_{3}^{2}h_{3}^{2}}} \ell = \sqrt[4]{\frac{3(1-0.3^{2})}{(60.007)^{2}(0.027)^{2}}} 0.5 = 0.505$$

Nonlinearity parameter:

$$\frac{a_1}{h_1} = \frac{60}{0.013} = 4615$$
 $\frac{p}{E} = \frac{35}{30 \times 10^6} = 1.17 \times 10^{-6}$

From figure 3, $n_1 = (1.012)(20.3) = 20.5$

It is now possible to find the stress factor by cross plotting. It is necessary to select stress factors for three nonlinearity parameters to determine accurately the

TABLE I. - STRESS FACTORS
FOR SAMPLE PROBLEM

Nonlinearity parameter,	Stress factor, K	Figure
1.0	1.51	6(f)
10	1.205	6(h)
100	1. 175	6(j)

stress factor for the specified example. For a thickness ratio of 0.929 and an overlap parameter of 0.505, stress factors are determined for the nonlinearity parameters in table I. Although these stress factors are for $\nu=1/3$ instead of $\nu=0.30$, it was noted previously (p. 6) that the result will be within ± 1.2 percent. A plot of the stress factors in figure 10 results in a stress factor of 1.19 for this example (n = 20.5). This same factor would be obtained for a corresponding unfilleted butt joint with about 30 percent mismatch and a corresponding filleted butt joint with about 34 percent mismatch.

CONCLUSIONS

Stress factors are presented in a set of design curves for circumferential joints in internally pressurized circular cylinders. Also presented are nonlinear edge influence coefficients for semi-infinite and short-length cylinders. Comparisons are made between linear and nonlinear results as well as between unfilleted butt, filleted butt, and overlap joints. Effects of stress concentrations at sharp corners are not considered. The following observations can be made:

- 1. For a change in thickness at a circumferential joint in a cylindrical pressure vessel, it is less detrimental if the thinner section is displaced inward rather than outward from the thicker section if mismatch is to occur.
- 2. If the thickness does not change at the junction of an unfilleted butt joint, a non-linear stress analysis is not required to determine the mismatch stresses.
- 3. For equal thicknesses at a filleted butt joint, the stress factor decreases as the nonlinearity increases. A decrease in the stress factor of up to 25 percent is possible for some filleted butt joints.

- 4. For overlap joints, a long overlap length is desirable only for nearly equal thicknesses. For widely differing thicknesses, a shorter length of overlap is preferable.
 - 5. For all three joints investigated, if the nonlinearity parameter

$$n = \frac{\sqrt[2]{3(1-\nu^2)}}{2} \frac{p}{E} \left(\frac{a}{h}\right)^2$$
 is about 0.8 or less, a difference between the linear and nonlinear results of no more than 5 percent is obtained.

Lewis Research Center,

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APPENDIX A

SYMBOLS

A	area of ring cross section, in. 2	Q	shear stress resultant, lb/in.						
a	middle surface radius of cylinder, in.	x	distance along meridian from discontinuity, in.						
b	width of ring $b_R + b_L$, in.	y	radial deflection of cylinder, in.						
c	thickness ratio, h_1/h_2	δ	edge-deflection influence coef-						
D	flexural rigidity of cylinder		ficient						
	$\frac{\text{Eh}^3}{12(1-\nu^2)}$, lb-in.	θ	meridional rotation, rad						
	·	λ	characteristic $\sqrt[4]{\frac{3(1-\nu^2)}{a^2h^2}}$, in. ⁻¹						
d	difference in middle surface radii $a_2 - a_1$, in.	ν	Poisson's ratio						
E	modulus of elasticity, psi	σ	normal stress, psi						
h	thickness of shell wall, in.	ω	edge-rotation influence coefficient						
I_{rr}	centroidal moment of inertia of	Subso	cripts:						
П	ring cross section about radial		discontinuity						
	axis, in. 4		effective						
K	stress factor (eq. (E9))	i	inner surface						
l 	length of overlap, in.	L	left of overlap cylinder or ring						
M	bending moment, (inlb)/in.	m	membrane ·						
m	percent mismatch, $\frac{a_2 - a_1}{(h_1 + h_2)/2} 100$		maximum						
			outer surface						
N	average axial stress resultant in	R	right of overlap cylinder or ring						
	wall of shell $\frac{p}{2} \frac{(a_1 + a_2)}{2}$, lb/in.	t	total						
	2 2	x	meridional direction						
n	nonlinearity parameter,	θ	circumferential direction						
	$\frac{\sqrt[4]{3(1-\nu^2)}}{2} \frac{p}{E} \left(\frac{a}{h}\right)^2$	1	thinner cylinder						
n	2 2 ()	2	thicker cylinder						
р	internal pressure, psi	3	overlap cylinder or ring						

APPENDIX B

SHEAR AND MOMENT AT UNFILLETED BUTT JOINT

The geometry and sign convention for an unfilleted butt joint is shown in figure 2(a). If the discontinuity forces are directed as shown in the figure, equilibrium of forces at the junction gives

$$Q_2 = -Q_1$$

$$M_2 = M_1 + Nd$$
(B1)

where

$$Nd = \left(\frac{p}{2} \frac{a_1 + a_2}{2}\right)(a_2 - a_1) = \frac{p}{4}(a_2^2 - a_1^2)$$
 (B2)

With the usual edge influence coefficient analysis, a simultaneous solution of the equations that express the equality of deflections and rotations at the joint and equation (B1) gives

$$Q_{1} = \frac{\left(\omega_{1, M_{1}} + \omega_{2, M_{2}}\right)\left[\operatorname{Nd\delta}_{2, M_{2}} - \operatorname{p}\left(\delta_{1, p} - \delta_{2, p}\right)\right] + \left(\delta_{1, M_{1}} - \delta_{2, M_{2}}\right)\left[\operatorname{Nd\omega}_{2, M_{2}} + \operatorname{p}\left(\omega_{1, p} + \omega_{2, p}\right)\right]}{\left(\delta_{1, Q_{1}} + \delta_{2, Q_{2}}\right)\left(\omega_{1, M_{1}} + \omega_{2, M_{2}}\right) - \left(\delta_{1, M_{1}} - \delta_{2, M_{2}}\right)\left(\omega_{1, Q_{1}} - \omega_{2, Q_{2}}\right)}$$

$$M_{1} = \frac{\left(\delta_{1, Q_{1}} + \delta_{2, Q_{2}}\right)\left[-\operatorname{Nd\omega}_{2, M_{2}} - \operatorname{p}\left(\omega_{1, p} + \omega_{2, p}\right)\right] - \left(\omega_{1, Q_{1}} - \omega_{2, Q_{2}}\right)\left[\operatorname{Nd\delta}_{2, M_{2}} - \operatorname{p}\left(\delta_{1, p} - \delta_{2, p}\right)\right]}{\left(\delta_{1, Q_{1}} + \delta_{2, Q_{2}}\right)\left(\omega_{1, M_{1}} + \omega_{2, M_{2}}\right) - \left(\delta_{1, M_{1}} - \delta_{2, M_{2}}\right)\left(\omega_{1, Q_{1}} - \omega_{2, Q_{2}}\right)}$$
(B3)

The well-known influence coefficients due to internal pressure for circular cylindrical shells are (ref. 15)

$$\delta_{1, p} = \frac{a_{1}^{2}}{Eh_{1}} \left(1 - \frac{\nu}{2} \right) \qquad \delta_{2, p} = \frac{a_{2}^{2}}{Eh_{2}} \left(1 - \frac{\nu}{2} \right)$$

$$\omega_{1, p} = 0 \qquad \qquad \omega_{2, p} = 0$$
(B4)

The remaining eight edge influence coefficients, which can be obtained from a reduction of equations given in reference 12, are

$$\begin{split} \delta_{1,\,Q_{1}} &= \frac{1}{2\lambda_{1}^{3}D_{1}} \left(\frac{\sqrt{1+n_{1}}}{1+2n_{1}} \right) & \delta_{2,\,Q_{2}} &= \frac{1}{2\lambda_{2}^{3}D_{2}} \left(\frac{\sqrt{1+n_{2}}}{1+2n_{2}} \right) \\ \delta_{1,\,M_{1}} &= \frac{1}{2\lambda_{1}^{2}D_{1}} \left(\frac{1}{1+2n_{1}} \right) & \delta_{2,\,M_{2}} &= \frac{1}{2\lambda_{2}^{2}D_{2}} \left(\frac{1}{1+2n_{2}} \right) \\ \omega_{1,\,Q_{1}} &= \frac{1}{2\lambda_{1}^{2}D_{1}} \left(\frac{1}{1+2n_{1}} \right) & \omega_{2,\,Q_{2}} &= \frac{1}{2\lambda_{2}^{2}D_{2}} \left(\frac{1}{1+2n_{2}} \right) \\ \omega_{1,\,M_{1}} &= \frac{1}{\lambda_{1}^{2}D_{1}} \left(\frac{\sqrt{1+n_{1}}}{1+2n_{1}} \right) & \omega_{2,\,M_{2}} &= \frac{1}{\lambda_{2}^{2}D_{2}} \left(\frac{\sqrt{1+n_{2}}}{1+2n_{2}} \right) \end{split} \end{split}$$

These are nonlinear edge influence coefficients, and when the nonlinearity parameter n is set equal to zero, the usual linear edge influence coefficients result.

Substitution of the edge influence coefficients of equations (B4) and (B5) in equation (B3) gives the discontinuity shear and moment for the thinner shell. The substitution of equation (B3) into equation (B1) gives the discontinuity shear and moment for the thicker shell. Once the discontinuity shear and moment have been determined, it is possible to find the deflection and stresses at any location (appendix E).

APPENDIX C

SHEAR AND MOMENT AT FILLETED BUTT JOINT

In the filleted butt joint, the structure is imagined cut at each discontinuity and separated into three components, with the center component treated as a ring of rigid cross section. The notation and sign convention used in this analysis is given in figure 2(b).

The combination of compatibility and equilibrium equations for each discontinuity gives four equations for determining the shear and moment on each cylindrical section:

$$g_{1} = Q_{1}(\delta_{3R,Q_{R}} - \delta_{1,Q_{1}}) + Q_{2}(\delta_{3R,Q_{L}}) + M_{1}(\delta_{3R,M_{R}} - \delta_{1,M_{1}}) + M_{2}(\delta_{3R,M_{L}})$$

$$g_{2} = Q_{1}(\delta_{3L,Q_{R}}) + Q_{2}(\delta_{3L,Q_{L}} - \delta_{2,Q_{2}}) + M_{1}(\delta_{3L,M_{R}}) + M_{2}(\delta_{3L,M_{L}} - \delta_{2,M_{2}})$$

$$g_{3} = Q_{1}(\omega_{3R,Q_{R}} - \omega_{1,Q_{1}}) + Q_{2}(\omega_{3R,Q_{L}}) + M_{1}(\omega_{3R,M_{R}} - \omega_{1,M_{1}}) + M_{2}(\omega_{3R,M_{L}})$$

$$g_{4} = Q_{1}(\omega_{3L,Q_{R}}) + Q_{2}(\omega_{3L,Q_{L}} - \omega_{2,Q_{2}}) + M_{1}(\omega_{3L,M_{R}}) + M_{2}(\omega_{3L,M_{L}} - \omega_{2,M_{2}})$$

$$(C1)$$

where

$$g_{1} = p(\delta_{1}, p - \delta_{3R}, p) + N \left[(a_{1} - a_{3})\delta_{3R}, M_{R} + (a_{2} - a_{3})\delta_{3R}, M_{L} \right]$$

$$g_{2} = p(\delta_{2}, p - \delta_{3L}, p) + N \left[(a_{1} - a_{3})\delta_{3L}, M_{R} + (a_{2} - a_{3})\delta_{3L}, M_{L} \right]$$

$$g_{3} = p(\omega_{1}, p - \omega_{3R}, p) + N \left[(a_{1} - a_{3})\omega_{3R}, M_{R} + (a_{2} - a_{3})\omega_{3R}, M_{L} \right]$$

$$g_{4} = p(\omega_{2}, p - \omega_{3L}, p) + N \left[(a_{1} - a_{3})\omega_{3L}, M_{R} + (a_{2} - a_{3})\omega_{3L}, M_{L} \right]$$
(C2)

For solution of these equations, 32 influence coefficients are required. The 20 influence coefficients for the ring section are determined by modification of the equations for a standard ring analysis, given in references 13 and 14. These influence coefficients are

$$\begin{split} \delta_{3L,Q_L} &= -\frac{a_3^2}{E} \left[\frac{1}{A} + \frac{b_L^2}{I_{rr} + nbh^2 / \sqrt{3(1 - \nu^2)}} \right] & \delta_{3R,Q_R} &= -\frac{a_3^2}{E} \left[\frac{1}{A} + \frac{b_R^2}{I_{rr} + nbh^2 / \sqrt{3(1 - \nu^2)}} \right] \\ \delta_{3L,Q_R} &= -\frac{a_3^2}{E} \left[\frac{1}{A} - \frac{b_R b_L}{I_{rr} + nbh^2 / \sqrt{3(1 - \nu^2)}} \right] & \delta_{3R,Q_L} &= -\frac{a_3^2}{E} \left[\frac{1}{A} - \frac{b_R b_L}{I_{rr} + nbh^2 / \sqrt{3(1 - \nu^2)}} \right] \\ \delta_{3L,M_L} &= \frac{a_3^2 b_L}{E \left[I_{rr} + nbh^2 / \sqrt{3(1 - \nu^2)} \right]} & \delta_{3R,M_R} &= \frac{a_3^2 b_R}{E \left[I_{rr} + nbh^2 / \sqrt{3(1 - \nu^2)} \right]} \\ \delta_{3L,M_R} &= -\frac{a_3^2 b_L}{E \left[I_{rr} + nbh^2 / \sqrt{3(1 - \nu^2)} \right]} & \delta_{3R,M_L} &= -\frac{a_3^2 b_R}{E \left[I_{rr} + nbh^2 / \sqrt{3(1 - \nu^2)} \right]} \\ \delta_{3L,Q_L} &= \frac{a_3^2 b_L}{E \left[I_{rr} + nbh^2 / \sqrt{3(1 - \nu^2)} \right]} & \omega_{3R,Q_R} &= \frac{a_3^2 b_R}{E \left[I_{rr} + nbh^2 / \sqrt{3(1 - \nu^2)} \right]} \\ \omega_{3L,Q_R} &= -\frac{a_3^2 b_R}{E \left[I_{rr} + nbh^2 / \sqrt{3(1 - \nu^2)} \right]} & \omega_{3R,Q_R} &= -\frac{a_3^2 b_L}{E \left[I_{rr} + nbh^2 / \sqrt{3(1 - \nu^2)} \right]} \\ \omega_{3L,M_L} &= -\frac{a_3^2}{E \left[I_{rr} + nbh^2 / \sqrt{3(1 - \nu^2)} \right]} & \omega_{3R,M_R} &= -\frac{a_3^2}{E \left[I_{rr} + nbh^2 / \sqrt{3(1 - \nu^2)} \right]} \\ \omega_{3L,M_R} &= -\frac{a_3^2}{E \left[I_{rr} + nbh^2 / \sqrt{3(1 - \nu^2)} \right]} & \omega_{3R,M_R} &= -\frac{a_3^2}{E \left[I_{rr} + nbh^2 / \sqrt{3(1 - \nu^2)} \right]} \\ \omega_{3L,M_R} &= -\frac{a_3^2}{E \left[I_{rr} + nbh^2 / \sqrt{3(1 - \nu^2)} \right]} & \omega_{3R,M_L} &= -\frac{a_3^2}{E \left[I_{rr} + nbh^2 / \sqrt{3(1 - \nu^2)} \right]} \\ \omega_{3L,M_R} &= -\frac{a_3^2}{E \left[I_{rr} + nbh^2 / \sqrt{3(1 - \nu^2)} \right]} & \omega_{3R,M_L} &= -\frac{a_3^2}{E \left[I_{rr} + nbh^2 / \sqrt{3(1 - \nu^2)} \right]} \\ \omega_{3L,M_R} &= -\frac{a_3^2}{E \left[I_{rr} + nbh^2 / \sqrt{3(1 - \nu^2)} \right]} & \omega_{3R,M_L} &= -\frac{a_3^2}{E \left[I_{rr} + nbh^2 / \sqrt{3(1 - \nu^2)} \right]} \\ \omega_{3L,M_R} &= -\frac{a_3^2}{E \left[I_{rr} + nbh^2 / \sqrt{3(1 - \nu^2)} \right]} & \omega_{3R,M_L} &= -\frac{a_3^2}{E \left[I_{rr} + nbh^2 / \sqrt{3(1 - \nu^2)} \right]} \\ \omega_{3L,M_R} &= -\frac{a_3^2}{E \left[I_{rr} + nbh^2 / \sqrt{3(1 - \nu^2)} \right]} & \omega_{3R,M_L} &= -\frac{a_3^2}{E \left[I_{rr} + nbh^2 / \sqrt{3(1 - \nu^2)} \right]} \\ \omega_{3L,M_R} &= -\frac{a_3^2}{E \left[I_{rr} + nbh^2 / \sqrt{3(1 - \nu^2)} \right]} & \omega_{3R,M_R} &= -\frac{a_3^2}{E \left[I_{rr} + nbh^2 / \sqrt{3(1 - \nu^2)} \right]} \\ \omega_{3L,M_R} &= -\frac{a_3^2}{E \left[I_{rr} + nbh^2 / \sqrt{3(1 - \nu^2)} \right]} & \omega_{3L,M_R} &= -\frac$$

 $\omega_{3L, p} = 0$

The 12 influence coefficients for the cylindrical components are given in equations (B4) and (B5).

In order to obtain the stress distribution for a filleted butt joint, it is first necessary to substitute the 32 influence coefficients in equations (C1) and (C2). Then, from a solution of these four equations, the unknown discontinuity shears and moments are determined. Once these are known, the deflection and stresses at any location of the cylindrical portions can be found (appendix E).

APPENDIX D

SHEAR AND MOMENT AT OVERLAP JOINT

The analysis at the overlap joint is very similar to the filleted butt joint analysis. In the overlap joint, the structure is imagined cut at each discontinuity and separated into three components, with the center component treated as a short-length cylinder. The notation and sign convention used in this analysis is given in figure 2(c). The combination of compatibility and equilibrium equations for each discontinuity gives four equations for determining the shear and moment on each cylindrical section (see eqs. (C1) and (C2)).

TABLE II. - CHANGES FOR OVERLAP EQUATIONS WHEN NONLINEARITY

PARAMETER $n_3 \le 1$

Terms in eqs. (D1) and (E5) for $n_3 > 1$	Changes necessary to use eqs. (D1) and (E5) when $n_3 \le 1$					
	n ₃ < 1	n ₃ = 1				
n ₃ - 1	1 - n ₃					
$\sinh(\lambda_3 \ell \sqrt{n_3 - 1})$	$\sin(\lambda_3 \ell \sqrt{1 - n_3})$					
$\cosh(\lambda_3 \ell \sqrt{n_3 - 1})$	$\cos\left(\lambda_3\ell\sqrt{1-n_3}\right)$					
n ₃		1				
$\frac{\sinh\left(\lambda_3\ell\sqrt{n_3-1}\right)}{\sqrt{n_3-1}}$		$\lambda_3\ell$				

The 16 edge influence coefficients due to edge shear force and bending moment for the overlap cylinder were determined from equations for a finite-length beam on an elastic foundation with free ends loaded by both an axial force and a unit-concentrated shear force or bending moment. The edge influence coefficients were taken from a modification of equations (104) and (105) of reference 12. These coefficients are dependent on whether the nonlinearity parameter n_3 is less than, equal to, or greater than 1. The influence coefficients for $n_3 > 1$ are

$$\begin{split} \delta_{3L,Q_L} &= \delta_{3R,Q_R} = -\frac{1}{2\lambda_D^2 O_3} \frac{\sqrt{1+n_3}}{1+2n_3} \frac{\sinh(\lambda_3 t \sqrt{1+n_3}) \cosh(\lambda_2 t \sqrt{1+n_3})}{\sinh^2(\lambda_3 t \sqrt{1+n_3}) - \sqrt{\frac{1+n_3}{n_3-1} \frac{1+2n_3}{1+2n_3}} \sinh(\lambda_3 t \sqrt{n_3-1}) \cosh(\lambda_2 t \sqrt{n_3-1})} \\ \delta_{3L,Q_R} &= \delta_{3R,Q_L} = -\frac{1}{2\lambda_D^2 O_3} \frac{\sqrt{1+n_3}}{1+2n_3} \frac{\sinh(\lambda_3 t \sqrt{1+n_3}) \cosh(\lambda_3 t \sqrt{1+n_3})}{\sinh^2(\lambda_3 t \sqrt{1+n_3}) - \frac{1+n_3}{n_3-1} \frac{1-2n_3}{1+2n_3} \sinh(\lambda_3 t \sqrt{n_3-1}) \cosh(\lambda_3 t \sqrt{1+n_3})} \\ \delta_{3L,M_L} &= \delta_{3R,M_R} = \frac{1}{2\lambda_D^2 O_3} \frac{1}{1+2n_3} \frac{\sinh(\lambda_3 t \sqrt{1+n_3}) + \frac{1+n_3}{n_3-1} \frac{1-2n_3}{1+2n_3} \sinh^2(\lambda_3 t \sqrt{1+n_3}) - \frac{1+n_3}{n_3-1} \frac{1-2n_3}{1+2n_3} \frac{\sinh^2(\lambda_3 t \sqrt{1+n_3}) - \frac{1+n_3}{n_3-1} \frac{1-2n_3}{1+2n_3} \sinh^2(\lambda_3 t \sqrt{1+n_3})}{\sinh^2(\lambda_3 t \sqrt{1+n_3}) - \frac{1+n_3}{n_3-1} \frac{1-2n_3}{1+2n_3} \sinh^2(\lambda_3 t \sqrt{n_3-1})} \\ \delta_{3L,M_R} &= \delta_{3R,M_L} = -\frac{1}{\lambda_D^2 O_3} \frac{1}{1+2n_3} \frac{\sinh^2(\lambda_3 t \sqrt{1+n_3}) - \frac{1+n_3}{n_3-1} \frac{1-2n_3}{1+2n_3} \sinh^2(\lambda_3 t \sqrt{n_3-1})}{\sinh^2(\lambda_3 t \sqrt{1+n_3}) - \frac{1+n_3}{n_3-1} \frac{1-2n_3}{1+2n_3} \sinh^2(\lambda_3 t \sqrt{n_3-1})} \\ \delta_{3L,Q_R} &= \delta_{3R,Q_R} = \frac{1}{2\lambda_D^2 O_3} \frac{1}{1+2n_3} \frac{\sinh^2(\lambda_3 t \sqrt{1+n_3}) - \frac{1+n_3}{n_3-1} \frac{1-2n_3}{1+2n_3} \sinh^2(\lambda_3 t \sqrt{n_3-1})}{\sinh^2(\lambda_3 t \sqrt{1+n_3}) - \frac{1+n_3}{n_3-1} \frac{1-2n_3}{1+2n_3} \sinh^2(\lambda_3 t \sqrt{n_3-1})} \\ \delta_{3L,M_R} &= \delta_{3R,Q_R} = \frac{1}{2\lambda_D^2 O_3} \frac{1}{1+2n_3} \frac{\sinh^2(\lambda_3 t \sqrt{1+n_3}) - \frac{1+n_3}{n_3-1} \frac{1-2n_3}{1+2n_3} \sinh^2(\lambda_3 t \sqrt{n_3-1})} \\ \delta_{3L,Q_R} &= \delta_{3R,Q_R} = \frac{1}{2\lambda_D^2 O_3} \frac{1}{1+2n_3} \frac{\sinh^2(\lambda_3 t \sqrt{1+n_3}) - \frac{1+n_3}{n_3-1} \frac{1-2n_3}{1+2n_3} \sinh^2(\lambda_3 t \sqrt{n_3-1})}{\sinh^2(\lambda_3 t \sqrt{1+n_3}) - \frac{1+n_3}{n_3-1} \frac{1-2n_3}{1+2n_3} \frac{1+2n_3}{2} \sinh^2(\lambda_3 t \sqrt{n_3-1})} \\ \delta_{3L,Q_R} &= \delta_{3R,Q_R} = \frac{1}{2\lambda_D^2 O_3} \frac{1}{1+2n_3} \frac{\sinh^2(\lambda_3 t \sqrt{1+n_3}) - \frac{1+n_3}{n_3-1} \frac{1-2n_3}{1+2n_3} \frac{1+2n_3}{2} \sinh^2(\lambda_3 t \sqrt{n_3-1})} \\ \delta_{3L,Q_R} &= \delta_{3R,Q_R} = \frac{1}{2\lambda_D^2 O_3} \frac{1}{1+2n_3} \frac{\sinh^2(\lambda_3 t \sqrt{1+n_3}) - \frac{1+n_3}{n_3-1} \frac{1-2n_3}{1+2n_3} \frac{1-2n_3}{2} \sinh^2(\lambda_3 t \sqrt{n_3-1})} \\ \delta_{3L,Q_R} &= \delta_{3R,Q_R} = \frac{1}{\lambda_D^2 O_3} \frac{1-2n_3}{1+2n_3} \frac{\sinh(\lambda_3 t \sqrt{1+n_3}) - \frac{1+n_3}{n_3-1} \frac{1-2n_3}{1+2n_3} \frac{1-2n_3}{2} \sinh^2(\lambda_3 t \sqrt{n_3-1})} \\ \delta_{3L,Q_R} &= \delta_{3R,Q_R} = \frac{1}{\lambda_D^2 O_3} \frac{1-2n_3}{1+2n_3} \frac$$

Because of complexity and length, the equations are presented only for $n_3 > 1$, where nonlinearity effects are more important. The changes necessary for use of these equations when $n_2 \le 1$, however, will be presented.

When $n_3 < 1$, certain terms in equation (D1) become imaginary numbers. When the two sets of equations are compared, it is found that three changes in equation (D1) will result in the equations for $n_3 < 1$. These changes are given in table II.

When $n_3 = 1$, an indeterminate quantity in equation (D1) becomes

$$\frac{\sinh\left(\lambda_3 \ell \sqrt{n_3 - 1}\right)}{\sqrt{n_3 - 1}} = \frac{0}{0} \stackrel{\mathbb{L}}{=} \lambda_3 \ell$$

after L'Hôpital's rule is applied. Therefore, equation (D1) will be applicable for $n_3 = 1$ if the changes given in table II are performed.

The four influence coefficients for component 3 (the short overlap cylinder), which are due to internal pressure, are determined in the same manner as for components 1 and 2. These can be stated as

$$\delta_{3R, p} = \delta_{3L, p} = \frac{a_3^2}{E(h_1 + h_2)} \left(1 - \frac{\nu}{2}\right)$$

$$\omega_{3R, p} = \omega_{3L, p} = 0$$
(D2)

The remaining 12 nonlinear influence coefficients are for the cylindrical components 1 and 2 and are given in equations (B4) and (B5).

In order to determine the discontinuity shear forces and bending moments for an overlap joint, it is first necessary to substitute the 32 influence coefficients in equations (C1) and (C2). From a solution of these four equations, the unknown discontinuity shear forces and bending moments are determined. Once these are known, the deflection and stresses at any location in the cylindrical portions can be found (appendix E).

APPENDIX E

STRESS DETERMINATION

The total stress at any location in the cylindrical components is composed of membrane stresses and discontinuity stresses. The well-known membrane stresses are as follows:

Meridional membrane stress:

$$\sigma_{x, m, i} = \sigma_{x, m, o} = \frac{pa}{2h}$$
 (E1)

Circumferential membrane stress:

$$\sigma_{\theta, m, i} = \sigma_{\theta, m, o} = \frac{pa}{h}$$
 (E2)

The discontinuity stresses are determined from the equations for a semi-infinite right circular cylindrical shell subjected to the action of shear forces and bending moments uniformly distributed along the edge x = 0. These discontinuity stresses are the following:

Meridional discontinuity stress:

$$\sigma_{x, d, i} = -\sigma_{x, d, o} = \frac{6}{h^2} M_{d, x}$$
 (E3)

Circumferential discontinuity stress:

$$\sigma_{\theta, d, i} = \frac{E}{a} y_{d, x} + \frac{6\nu}{h^2} M_{d, x}$$

$$\sigma_{\theta, d, o} = \frac{E}{a} y_{d, x} - \frac{6\nu}{h^2} M_{d, x}$$
(E4)

The discontinuity deflection and bending moment distribution as a function of x can be determined by combining equations (100) and (101) of reference 12. These equations are dependent on whether the nonlinearity parameter n is less than, equal to, or greater than 1. For brevity, the equations are presented only for n > 1 where nonlinear effects

are more important. For $n \le 1$, the equations should be changed in the manner given in table II. The equations for n > 1 are

$$\begin{split} y_{d,\,x} = & \frac{\exp\left(-\lambda x\,\sqrt{1+n}\right)}{2\lambda^2 D(1+2n)} \left\{ M \left[\cosh\left(\lambda x\,\sqrt{n-1}\right) - \sqrt{\frac{1+n}{n-1}}\,\sinh\left(\lambda x\,\sqrt{n-1}\right) \right] \right. \\ & \left. + \frac{Q}{\lambda} \left[\sqrt{1+n}\,\cosh\left(\lambda x\,\sqrt{n-1}\right) + \frac{n}{\sqrt{n-1}}\,\sinh\left(\lambda x\,\sqrt{n-1}\right) \right] \right\} \\ M_{d,\,x} = & \exp\left(-\lambda x\,\sqrt{1+n}\right) \left\{ M \left[\cosh\left(\lambda x\,\sqrt{n-1}\right) + \sqrt{\frac{1+n}{n-1}} \left(\frac{1-2n}{1+2n}\right) \sinh\left(\lambda x\,\sqrt{n-1}\right) \right] \right. \\ & \left. + \frac{Q}{\lambda(1+2n)} \frac{\sinh\left(\lambda x\,\sqrt{n-1}\right)}{\sqrt{n-1}} \right\} \end{split}$$

The total stress is the sum of the membrane and discontinuity stresses. The addition of corresponding stresses gives

$$\sigma_{\mathbf{x}, \mathbf{t}, \mathbf{i}} = \sigma_{\mathbf{x}, \mathbf{m}, \mathbf{i}} + \sigma_{\mathbf{x}, \mathbf{d}, \mathbf{i}}$$

$$\sigma_{\mathbf{x}, \mathbf{t}, \mathbf{o}} = \sigma_{\mathbf{x}, \mathbf{m}, \mathbf{o}} + \sigma_{\mathbf{x}, \mathbf{d}, \mathbf{o}}$$

$$\sigma_{\theta, \mathbf{t}, \mathbf{i}} = \sigma_{\theta, \mathbf{m}, \mathbf{i}} + \sigma_{\theta, \mathbf{d}, \mathbf{i}}$$

$$\sigma_{\theta, \mathbf{t}, \mathbf{o}} = \sigma_{\theta, \mathbf{m}, \mathbf{o}} + \sigma_{\theta, \mathbf{d}, \mathbf{o}}$$

$$(E6)$$

The effective total stresses, defined by using the energy of distortion theory for a biaxial state of stress condition, are

$$\sigma_{e, t, i} = \sqrt{(\sigma_{x, t, i})^{2} + (\sigma_{\theta, t, i})^{2} - (\sigma_{x, t, i})(\sigma_{\theta, t, i})}$$

$$\sigma_{e, t, o} = \sqrt{(\sigma_{x, t, o})^{2} + (\sigma_{\theta, t, o})^{2} - (\sigma_{x, t, o})(\sigma_{\theta, t, o})}$$
(E7)

The maximum effective membrane stress, which is always in the thinner shell, is

$$\sigma_{e, m, 1} = \sqrt{(\sigma_{x, m, 1})^2 + (\sigma_{\theta, m, 1})^2} - (\sigma_{x, m, 1})(\sigma_{\theta, m, 1}) = \frac{\sqrt{3}}{2} \frac{pa_1}{h_1}$$
 (E8)

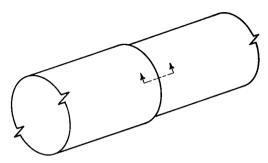
The stress factor K is the ratio of the largest value given by equation (E7) to equation (E8) and is

$$K = \frac{\sigma_{e, t, max}}{\sigma_{e, m, 1}}$$
 (E9)

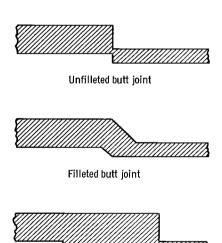
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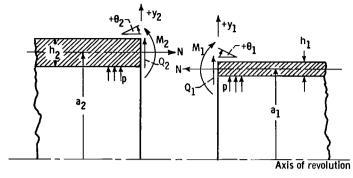
(a) Junction of two semi-infinite circular cylinders.



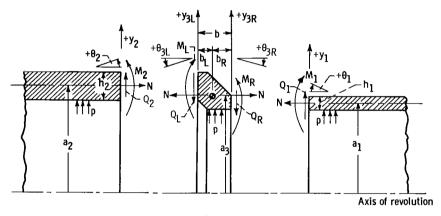
(b) Illustration of joints.

Figure 1. - Configurations investigated.

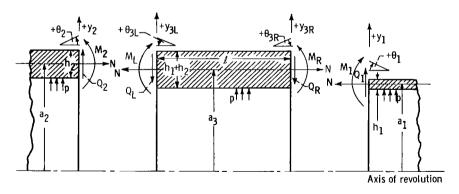
Overlap joint



(a) Unfilleted butt joint.



(b) 45° Filleted butt joint.



(c) Overlap joint

Figure 2. - Geometry and sign convention.

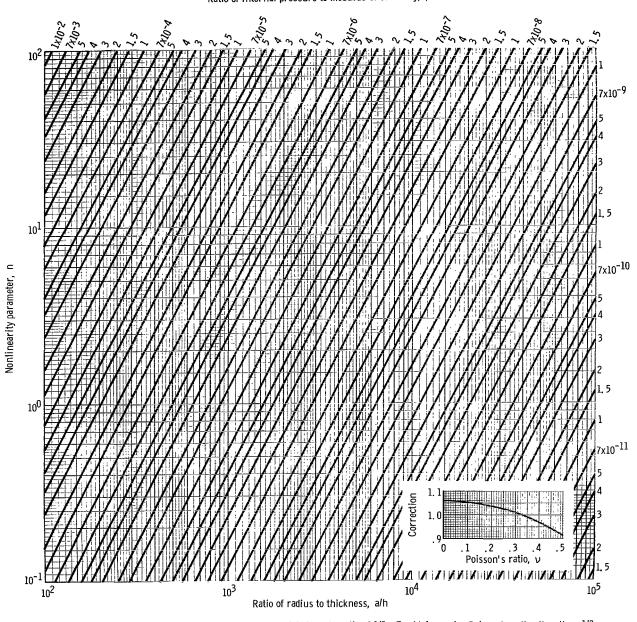
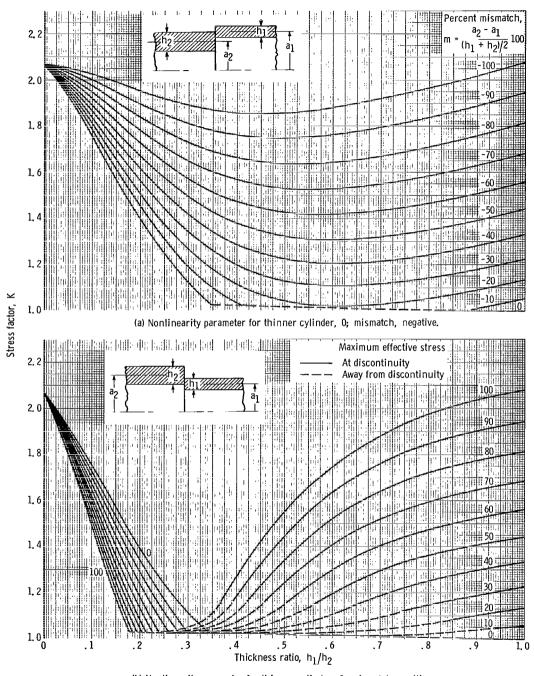
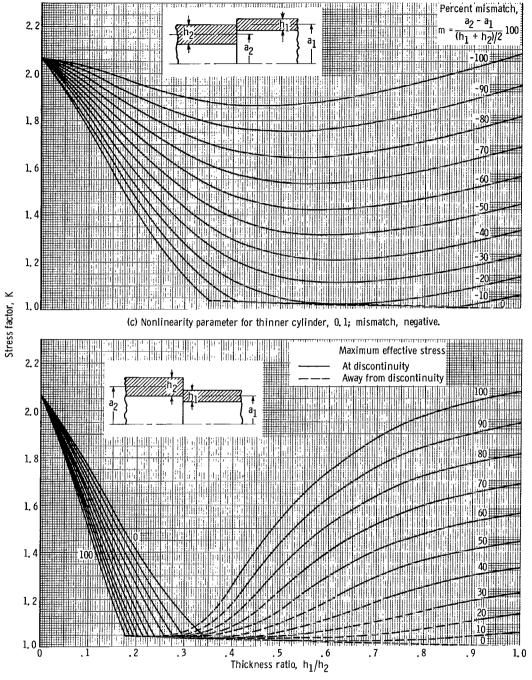


Figure 3. – Determination of nonlinearity parameter $\,$ n at Poisson's ratio of 1/3. To obtain $\,$ n for Poisson's ratio other than 1/3, multiply value from figure by correction given.



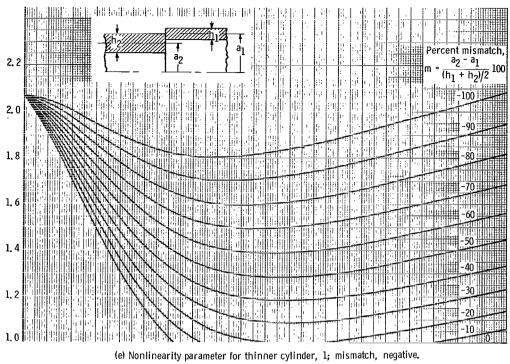
(b) Nonlinearity parameter for thinner cylinder, 0; mismatch, positive.

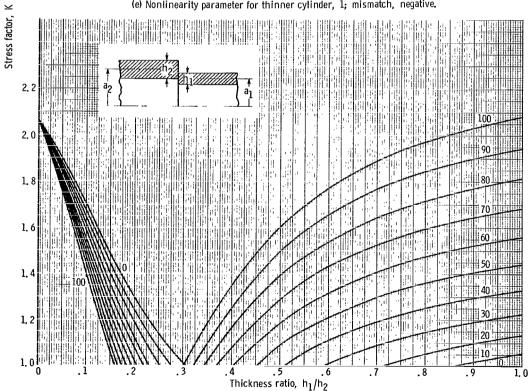
Figure 4. - Stress factors for unfilleted butt joint. Poisson's ratio, 1/3.



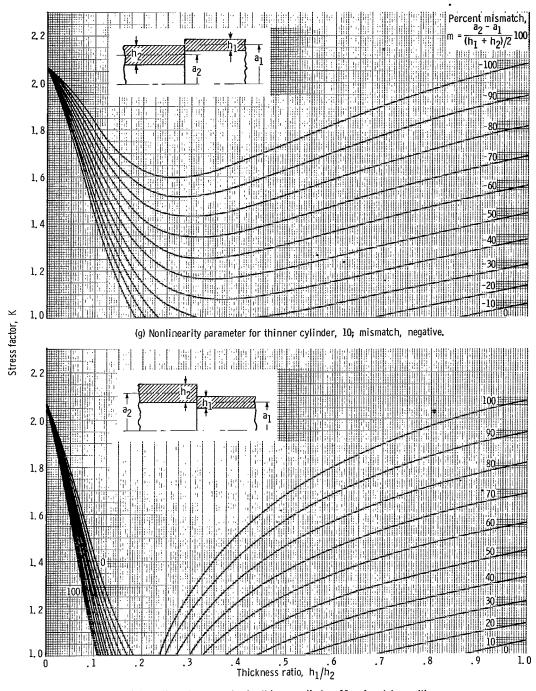
(d) Nonlinearity parameter for thinner cylinder, 0.1; mismatch, positive.

Figure 4. - Continued.

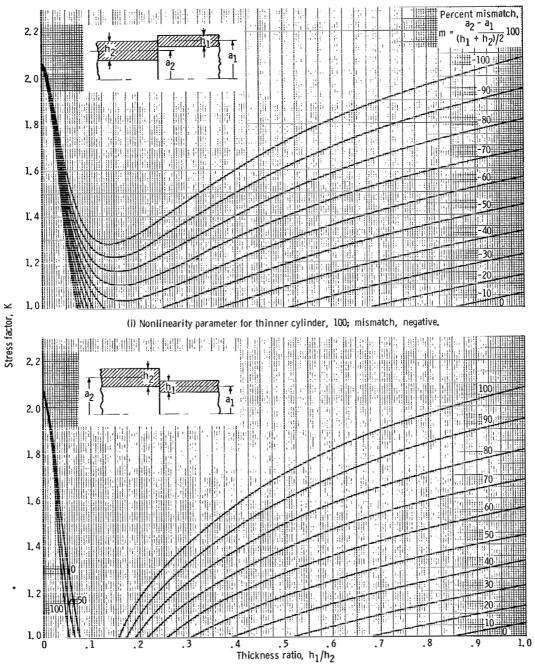




(f) Nonlinearity parameter for thinner cylinder, 1; mismatch, positive. Figure 4. - Continued.

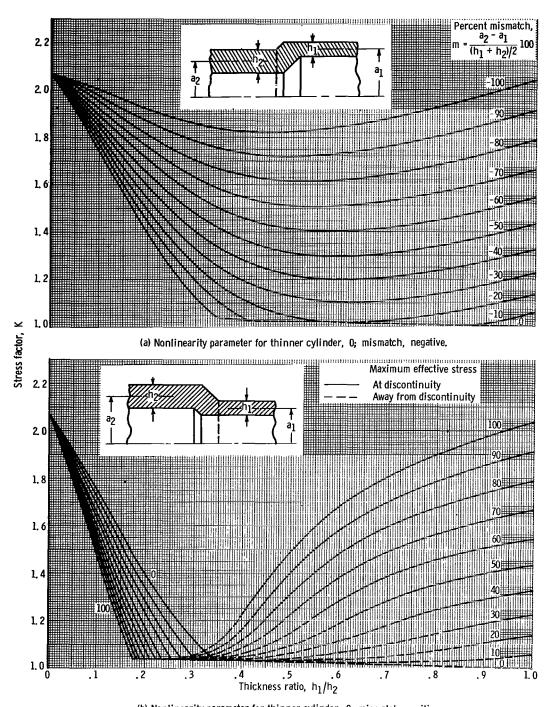


(h) Nonlinearity parameter for thinner cylinder, 10; mismatch, positive. Figure 4. - Continued.

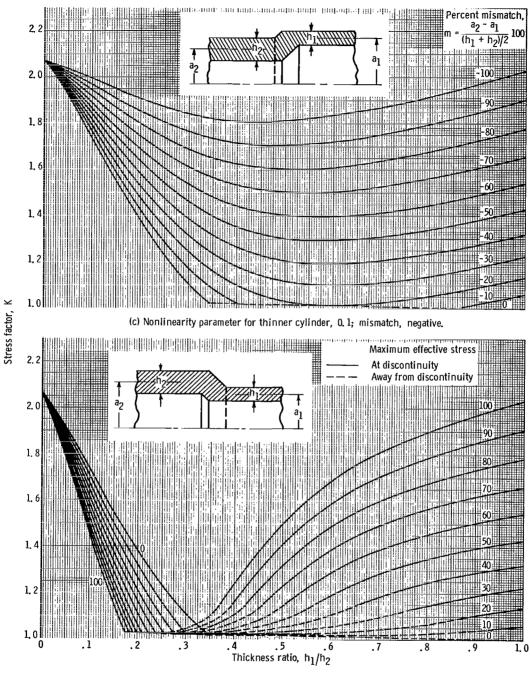


(j) Nonlinearity parameter for thinner cylinder, 100; mismatch, positive.

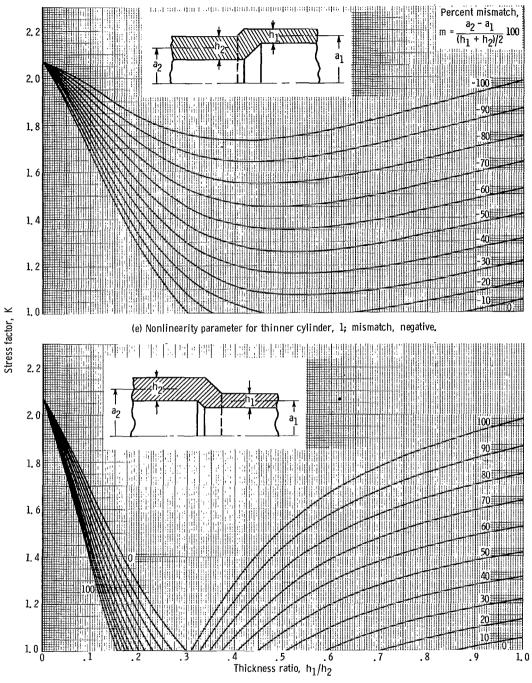
Figure 4. - Concluded.



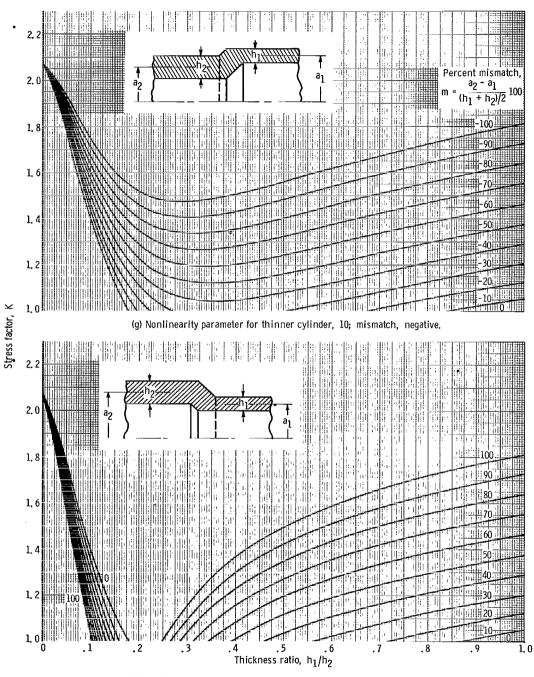
(b) Nonlinearity parameter for thinner cylinder, 0; mismatch, positive. Figure 5. - Stress factors for 45° filleted butt joint. Poisson's ratio, 1/3.



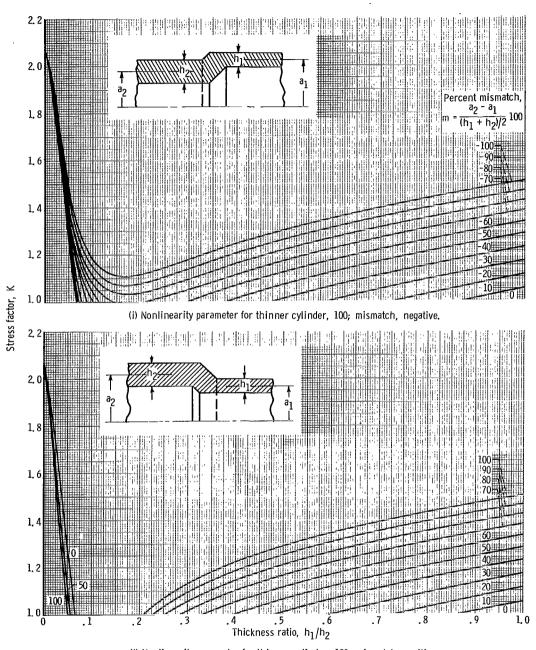
(d) Nonlinearity parameter for thinner cylinder, 0, 1; mismatch, positive. Figure 5. - Continued.



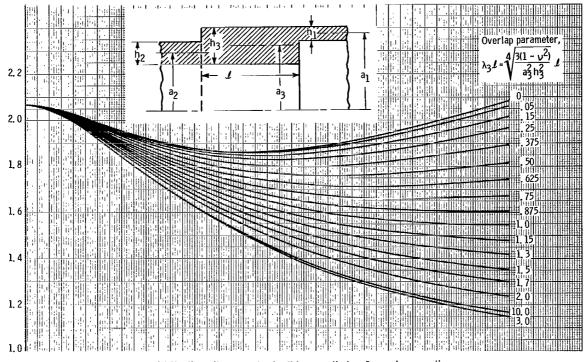
(f) Nonlinearity parameter for thinner cylinder, 1; mismatch, positive. $\mbox{Figure 5.} \ - \mbox{Continued}.$

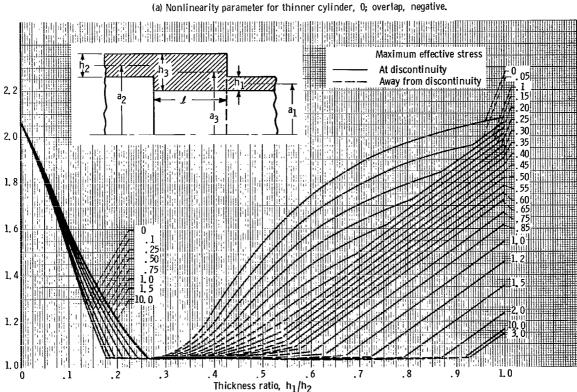


(h) Nonlinearity parameter for thinner cylinder, 10; mismatch, positive. Figure 5. - Continued.



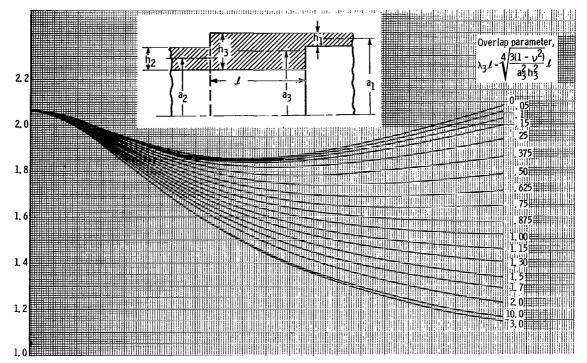
(j) Nonlinearity parameter for thinner cylinder, 100; mismatch, positive. Figure 5. – Concluded.



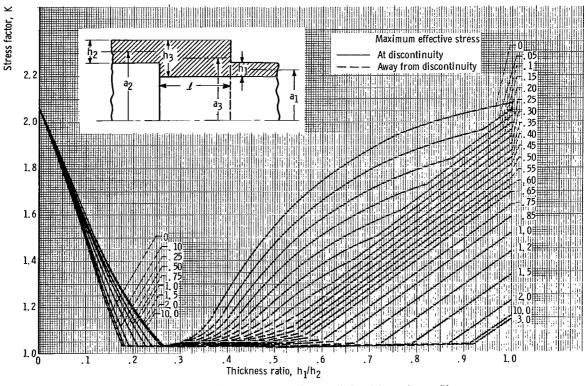


Stress factor, K

(b) Nonlinearity parameter for thinner cylinder, 0; overlap, positive. Figure 6. - Stress factors for overlap joint. Poisson's ratio, 1/3.

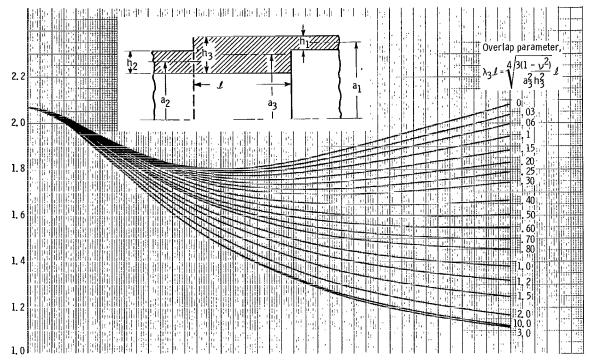


(c) Nonlinearity parameter for thinner cylinder. 0.1: overlap, negative.

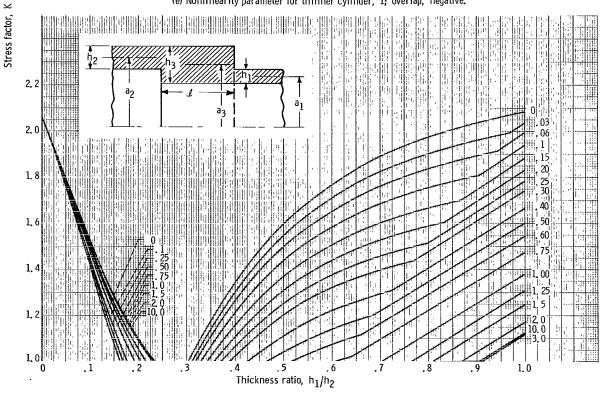


(d) Nonlinearity parameter for thinner cylinder, 0.1, overlap, positive.

Figure 6. - Continued.

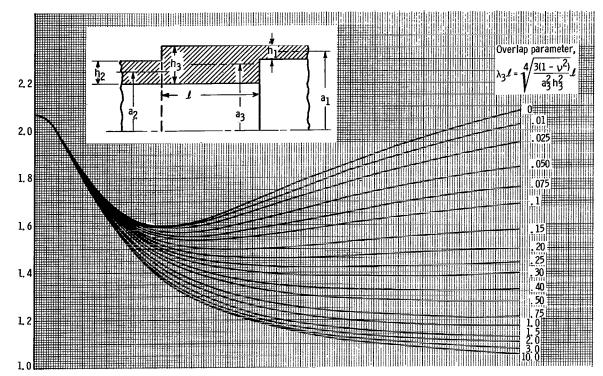




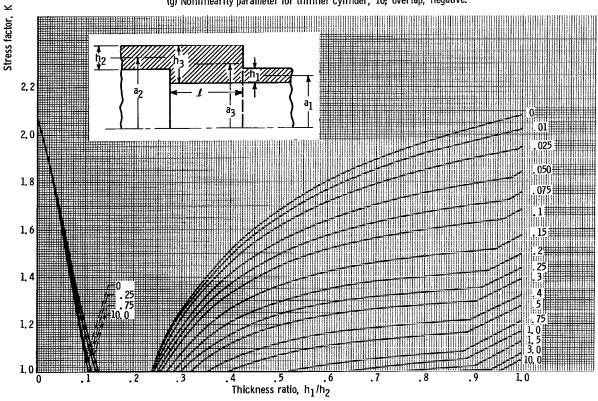


(f) Nonlinearity parameter for thinner cylinder, 1; overlap, positive.

Figure 6. - Continued.

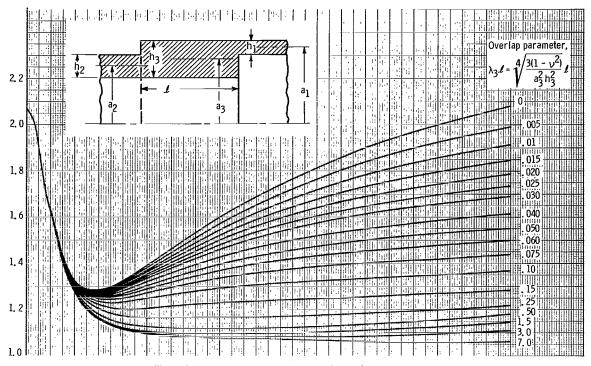


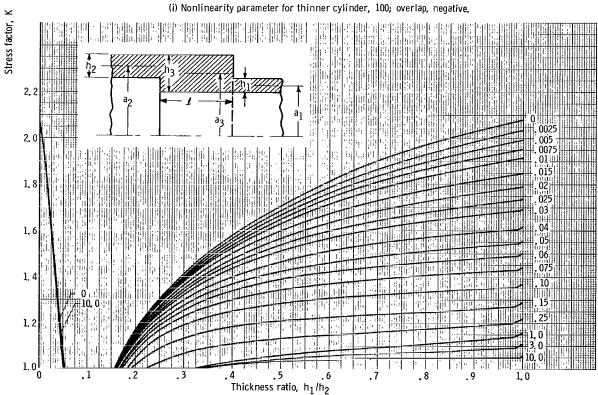




(h) Nonlinearity parameter for thinner cylinder, 10; overlap, positive.

Figure 6. - Continued.





(j) Nonlinearity parameter for thinner cylinder, 100; overlap, positive.Figure 6. - Concluded.

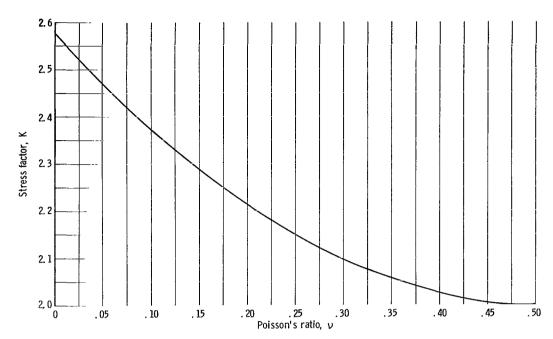


Figure 7. - Stress factor as function of Poisson's ratio for pressurized circular cylinder with fixed ends (longitudinal expansion permitted).

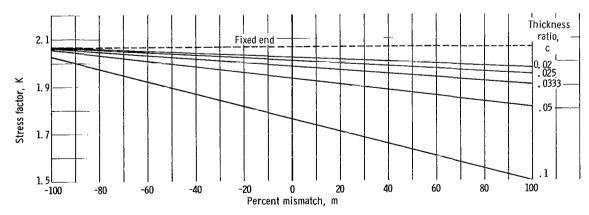


Figure 8. - Stress factors for unfilleted butt joints with small thickness ratios,

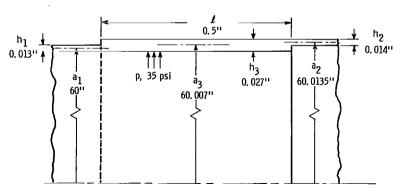


Figure 9. - Geometry and loading for sample problem. Modulus of elasticity, $30x10^6$ pounds per square inch; Poisson's ratio, 0.3.

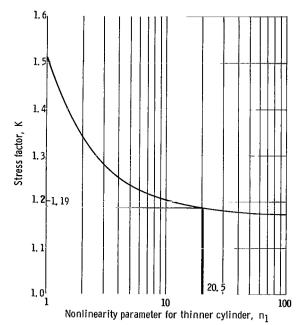


Figure 10. - Graphical determination of stress factor for sample problem.

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